## Lecture : Matrix Decompositions

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## 1 Symmetric Matrices

### 1.1 Symmetric and Hermitian Matrices

A symmetric matrix is a real matrix $S \in \mathbb{R}^{n \times n}$ such that $S=S^{T}$. A Hermitian matrix $H \in \mathbb{C}^{n \times n}$ is a matrix such that $H=H^{*}$. Symmetric matrices have all real eigenvalues and can be diagonalized by rotation matrices, ie. for every symmetric matrix $S$, there exists a rotation matrix $R$ and a diagonal matrix of real eigenvalues $D$ such that

$$
S=R D R^{T}=\left[\begin{array}{ccc}
\mid & & \mid  \tag{1}\\
r_{1} & \cdots & r_{n} \\
\mid & & \mid
\end{array}\right]\left[\begin{array}{ccc}
\lambda_{1} & & 0 \\
\vdots & \ddots & \vdots \\
0 & & \lambda_{n}
\end{array}\right]\left[\begin{array}{ccc}
- & r_{1}^{T} & - \\
& \vdots & \\
- & r_{n}^{T} & -
\end{array}\right]
$$

This means that multiplying by an $n \times n$ symmetric matrix corresponds to stretching in $n$ orthogonal coordinate directions. Symmetric matrices can be thought of as defining level sets of the quadratic form $h(x)=\frac{1}{2} x^{T} S x$. The vector field $\dot{x}=S x$ is then a gradient field, $\dot{x}=\frac{\partial h^{T}}{\partial x}$. Intuitively, $\dot{x}=S x$ can be thought of as flowing up a surface defined by $h(x)$. This is also called a conservative vector field in physics and the function $h(x)$ is typically related to the energy of a system (or some analog). The general condition for a nonlinear vector field $\dot{x}=f(x)$ to be conservative, ie. that $f(x)={\frac{\partial h^{T}}{\partial x}}^{\text {for }}$ fome $h(x)$ is given by $\frac{\partial f_{i}}{\partial x_{j}}=\frac{\partial f_{j}}{\partial x_{i}}$ for all $i, j$. This condition is derived from the fact that if such and $h(x)$ exists, then $\frac{\partial^{2} h}{\partial x_{i} \partial x_{j}}=\frac{\partial^{2} h}{\partial x_{j} \partial x_{i}}$.

Symmetric Matrices/Conservative Vector Field



### 1.2 Positive Definiteness

We say a symmetric matrix $S$ is positive definite if

$$
\begin{equation*}
x^{T} S x>0, \quad \text { for all } x \in \mathbb{R}^{n} \tag{2}
\end{equation*}
$$

Since we could pick $x$ to be any $r_{i}$, it follows that for a positive definite matrix, $\lambda_{i}>0$ for all $i$. If the "greater than" signs above are replaced with "greater than or equal" signs, ie. " $>$ " is replaced with " $\geq$ " then we say the matrix is positive semi-definite. Similarly, if we replace the "greater than" signs with "less than" signs we say that the matrix is negative definite or negative semi-definite.

### 1.3 Polar Decomposition

For any matrix $A \in \mathbb{R}^{m \times n}$, there are two positive semi-definite matrices $P=\left(A^{T} A\right)^{\frac{1}{2}}$ and $P^{\prime}=$ $\left(A A^{T}\right)^{\frac{1}{2}}$ (where the positive square root of each eigenvalue is taken) that are closely related to the "shape" of the matrix $A$. Similarly to how the magnitude of a complex number is defined by $|z|=\sqrt{z^{*} z}$, we can say that the "magnitude and shape" of $A$ is defined either by $\left(A^{T} A\right)^{\frac{1}{2}}$ or $\left(A A^{T}\right)^{\frac{1}{2}}$. Expanding out $A^{T} A$, we get

$$
\begin{aligned}
A^{T} A & =\left[\begin{array}{ccc}
- & A_{1}^{T} & - \\
& \vdots & \\
- & A_{n}^{T} & -
\end{array}\right]\left[\begin{array}{ccc}
\mid & & \mid \\
A_{1} & \cdots & A_{n} \\
\mid & & \mid
\end{array}\right] \\
& =\left[\begin{array}{ccc}
A_{1}^{T} A_{1} & \cdots & A_{1}^{T} A_{n} \\
\vdots & & \vdots \\
A_{n}^{T} A_{1} & \cdots & A_{n}^{T} A_{n}
\end{array}\right]=\left[\begin{array}{ccc}
\left|A_{1}\right|\left|A_{1}\right| \cos \left(\theta_{11}\right) & \cdots & \left|A_{1}\right|\left|A_{n}\right| \cos \left(\theta_{1 n}\right) \\
\vdots & & \vdots \\
\left|A_{n}\right|\left|A_{1}\right| \cos \left(\theta_{n 1}\right) & \cdots & \left|A_{n}\right|\left|A_{n}\right| \cos \left(\theta_{n n}\right)
\end{array}\right]
\end{aligned}
$$

We note that this matrix is fully determined by the size and relative orientation of the columns of $A$. Another way to say this is that applying an orthonormal transformation to all the columns of $A$ does not change $A^{T} A$. Indeed $(R A)^{T}(R A)=A^{T} R^{T} R A=A^{T} A$. Similarly the size and relative orientation of the rows of $A$ full determines $A A^{T}$. We can make precise the sense in which $A$ has the same shape as $P=\left(A^{T} A\right)^{\frac{1}{2}}$ by noting that $P$ and $A$ differ by a orthonormal transformation. Indeed,

$$
A=\underbrace{A\left(A^{T} A\right)^{-\frac{1}{2}}}_{R} \underbrace{\left(A^{T} A\right)^{\frac{1}{2}}}_{P}
$$

We note that we can check that $R^{T} R=I$

$$
\begin{aligned}
R^{T} R & =\left(A^{T} A\right)^{-\frac{1}{2}} A^{T} A\left(A^{T} A\right)^{-\frac{1}{2}} \\
& =\underbrace{\left(A^{T} A\right)^{-\frac{1}{2}}\left(A^{T} A\right)^{\frac{1}{2}}}_{I} \underbrace{\left(A^{T} A\right)^{\frac{1}{2}}\left(A^{T} A\right)^{-\frac{1}{2}}}_{I}=I
\end{aligned}
$$

and thus we have that rotating (and possibly reflecting) all the columns of the positive semidefinite matrix $P \succeq 0$ by one rotation gives $A$. Similarly a complex number $z=|z| e^{i \theta}$ can be created by starting with it's norm $|z| \geq 0$ and rotating it in the complex plane by $e^{i \theta}$. Similarly $A$ can be created from $\left(A A^{T}\right)^{\frac{1}{2}}$ by applying the orthonormal transformation $R^{\prime}=\left(A A^{T}\right)^{-\frac{1}{2}} A$.

$$
A=\underbrace{\left(A A^{T}\right)^{\frac{1}{2}}}_{P^{\prime}} \underbrace{\left(A A^{T}\right)^{-\frac{1}{2}} A}_{R^{\prime}}
$$

One can check that in fact for a square matrix, $R=R^{\prime}$.
This leads us to the polar decomposition. A square, invertible $A \in \mathbb{R}^{n \times n}$ can be written in a polar form similar to the polar decomposition of a complex number $z=\sqrt{z^{*} z} e^{i \theta}$. The

$$
A=R P=P^{\prime} R
$$

where

$$
\begin{aligned}
& P=\left(A^{*} A\right)^{\frac{1}{2}}=V \Sigma V^{*} \\
& P=\left(A A^{*}\right)^{\frac{1}{2}}=U \Sigma U^{*} \\
& R=\left(A A^{*}\right)^{\frac{-1}{2}} A=A\left(A^{*} A\right)^{\frac{-1}{2}}=U V^{*}
\end{aligned}
$$

Note that $P \succeq 0$ and $P^{\prime} \succeq 0$. Note also the There are two separate versions of the polar decomposition one with $P$ and one with $P^{\prime}$. Note also the connections between the polar decomposition and the singular value decomposition (see below). These relationships can be checked directly.

## 2 Singular Value Decomposition

The singular value decomposition (SVD) provides even more insight beyond the polar decomposition. The SVD is very general and can apply to any matrix $A \in \mathbb{C}^{m \times n}$ even if the matrix is not-invertible, not diagonalizable, or even not square or full-rank. As a result, it is an often used, powerful analysis tool. We will perform the derivation below for $A \in \mathbb{R}^{m \times n}$ but we note that the same derivation works for complex matrices using conjugate transposes.

To construct the SVD, we will assume that $A$ is fat or square and analyze $A^{T} A$. (A similar construction can be done if $A$ is tall using $A A^{T}$.). We first diagonalize $A^{T} A$ as

$$
A^{T} A=V D V^{T}
$$

where $V \in \mathbb{R}^{n \times n}$ is orthonormal, ie. $V^{T} V=I$, and $D \succeq 0$ is diagonal, real and positive semidefinite. We know this is possible since $A^{T} A$ is symmetric and thus has all real eigenvalues and orthogonal eigenvectors. We will also assume that the diagonal of $D$ is arranged in descending
order from the largest eigenvalue to the smallest. Since $A$ is fat or square, some portion of the diagonal of $D$ will be 0 's. We can take the positive square root of $D$ to get

$$
D^{\frac{1}{2}}=\left[\begin{array}{ll}
\Sigma & 0 \\
0 & 0
\end{array}\right]
$$

where

$$
\Sigma=\left[\begin{array}{ccc}
\sigma_{1} & \cdots & 0 \\
\vdots & & \vdots \\
0 & \cdots & \sigma_{k}
\end{array}\right]
$$

and the block zeros have the appropriate sizes. $\left\{\sigma_{i}\right\}_{i=1}^{k}$ are called singular values and are the positive square roots of the nonzero eigenvalues of $A^{T} A$. We can enumerate $V$ as

$$
V=\left[\begin{array}{ll}
V^{1} & V^{2}
\end{array}\right]
$$

where

$$
V^{1}=\left[\begin{array}{ccc}
\mid & & \mid \\
V_{1} & \cdots & V_{k} \\
\mid & & \mid
\end{array}\right], \quad V^{2}=\left[\begin{array}{ccc}
\mid & & \mid \\
V_{k+1} & \cdots & V_{n} \\
\mid & & \mid
\end{array}\right]
$$

where the columns of $V^{1}$ correspond to positive, non-zero singular values and the columns of $V^{2}$ are the eigenvectors for the zero eigenvalues of $A^{T} A$. We note that the columns of $V^{2}$ can be chosen somewhat arbitrarily as long as they are orthonormal and span $\mathcal{N}\left(A^{T} A\right)$. Since $\mathcal{N}\left(A^{T} A\right)=\mathcal{N}(A)$, $\mathcal{R}\left(V^{2}\right) \in \mathcal{N}(A)$ as well.

We note that any columns corresponding to repeated eigenvalues of $A^{T} A$ (including repeated zeros), may be arbitrarily chosen (as long as they are orthonormal). We can now define a matrix $U \in \mathbb{R}^{m \times m}$ as

$$
U=\left[\begin{array}{ll}
U^{1} & U^{2}
\end{array}\right]
$$

where

$$
U^{1}=\left[\begin{array}{ccc}
\mid & & \mid \\
U_{1} & \cdots & U_{k} \\
\mid & & \mid
\end{array}\right]=\left[\begin{array}{ccc}
\mid & & \mid \\
\frac{A V_{1}}{\sigma_{1}} & \cdots & \frac{A V_{k}}{\sigma_{k}} \\
\mid & & \mid
\end{array}\right] \quad U^{2}=\left[\begin{array}{ccc}
\mid & & \mid \\
U_{k+1} & \cdots & U_{m} \\
\mid & & \mid
\end{array}\right]
$$

Note that the columns of $U^{1}$ are orthonormal.

$$
\frac{V_{i}^{T} A^{T}}{\sigma_{i}} \frac{A V_{i}}{\sigma_{i}}=\frac{\sigma_{i}^{2}}{\sigma_{i}^{2}} V_{i}^{T} V_{i}=1, \quad \frac{V_{i}^{T} A^{T}}{\sigma_{i}} \frac{A V_{j}}{\sigma_{j}}=\frac{\sigma_{i}^{2}}{\sigma_{j}^{2}} V_{i}^{T} V_{j}=0
$$

The columns of $U_{2}$ can be chosen to complete an orthonormal basis for $\mathbb{R}^{m}$ and thus $U^{T} U=I$. We can then write

$$
U^{1} \Sigma=A V^{1}
$$

By adding $(0) U^{2}$ to the left-hand side, and since $V^{2} \in \mathcal{N}(A)$, we can write

$$
\begin{aligned}
& {\left[\begin{array}{ll}
U^{1} \Sigma+(0) U^{2} & 0
\end{array}\right]=A\left[V^{1} V^{2}\right]} \\
& {\left[\begin{array}{ll}
U^{1} & U^{2}
\end{array}\right]\left[\begin{array}{ll}
\Sigma & 0 \\
0 & 0
\end{array}\right]=A\left[V^{1} V^{2}\right]}
\end{aligned}
$$

right multiplying by $V^{T}$ gives

$$
\begin{aligned}
A & =U\left[\begin{array}{ll}
\Sigma & 0 \\
0 & 0
\end{array}\right] V^{T} \\
& =\left[\begin{array}{ll}
U^{1} & U^{2}
\end{array}\right]\left[\begin{array}{ll}
\Sigma & 0 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
V^{1 T} \\
V^{2 T}
\end{array}\right]
\end{aligned}
$$

This is the singular value decomposition. Note the following relationships

$$
\begin{array}{cc}
\mathcal{R}\left(U^{1}\right)=\mathcal{R}(A), & \mathcal{R}\left(V^{1}\right)=\mathcal{R}\left(A^{T}\right) \\
\mathcal{R}\left(U^{2}\right)=\mathcal{N}\left(A^{T}\right), & \mathcal{R}\left(V^{2}\right)=\mathcal{N}(A)
\end{array}
$$

Note also that the columns of $U$ are orthonormal eigenvectors of $A A^{T}$ with the singular values squared as eigenvalues. Indeed,

$$
A A^{T} U_{i}=\frac{A A^{T} A V_{i}}{\sigma_{i}}=\frac{\sigma_{i}^{2}}{\sigma_{i}} A V_{i}=\left(\sigma_{i}\right)^{2} U_{i}
$$

Thus this construction would have worked using $A A^{T}$ instead of $A^{T} A$ if desired.

### 2.1 Symmetric-Skew Symmetric/Helmholtz Decomposition

A square matrix $A \in \mathbb{R}^{n \times n}$ can be decomposed as follows

$$
A=\underbrace{\frac{1}{2}\left(A+A^{T}\right)}_{S}+\underbrace{\frac{1}{2}\left(A-A^{T}\right)}_{K}
$$

Note that $S=S^{T}$ is symmetric and $K=-K^{T}$ is skew-symmetric. This decomposition says that the space of real matrices is actually the direct sum of the space of symmetric matrices and the space of skew-symmetric matrices. Under the vectorized matrix inner product $\langle\cdot, \cdot\rangle=\operatorname{Tr}\left((\cdot)^{T}(\cdot)\right)$, we have that

$$
\begin{aligned}
\langle S, K\rangle & =\operatorname{Tr}\left(S^{T} K\right)=\sum_{i, j} S_{i j} K_{i j} \\
& =\sum_{i<j} S_{i j} K_{i j}+\sum_{i>j} S_{i j} K_{i j}+\sum_{i} S_{i i} K_{i i} \\
& =\sum_{i<j} S_{i j} K_{i j}+\sum_{j>i} S_{j i} K_{j i}+\sum_{i} S_{i i} K_{i i} \\
& =\sum_{i<j} S_{i j}\left(K_{i j}+K_{j i}\right)+\sum_{i} S_{i i} K_{i i}=0
\end{aligned}
$$

If the definition of positive definite is extended to non-symmetric matrices we have that

$$
x^{T} A x=x^{T} S x+x^{T} K x=x^{T} S x
$$

and thus $A$ is positive definite if and only if $S=\frac{1}{2}\left(A+A^{T}\right)$ is positive definite. In the context of vector fields, this means that any linear vector field $\dot{x}=A x$ can be decomposed into a conservative piece and a rotational piece.

$$
\dot{x}=A x=\underbrace{S x}_{\text {conservative }}+\underbrace{K x}_{\text {rotational }}
$$

which is a special application of the Helmholtz decomposition to linear vector fields.

## 3 Matrix vs. Complex Number Analogies

The polar decomposition and the symmetric/skew-symmetric decomposition provide several deep analogies between matrices and complex numbers. Symmetric matrices act a lot like real numbers and positive definite symmetric matrices act like positive real numbers. Among other things, they have real and positive real eigenvalues respectively. Skew-symmetric matrices act a lot like purely imaginary numbers, encoding rotational flow and having purely imaginary eigenvalues. We can even draw a "complex plane" of sorts for matrices with symmetric matrices as the real axis and skew-symmetric matrices as the vertical axis. As shown above, the fact that these two subspaces are orthogonal is accurate. The Cartesian description of complex numbers is analogous to the symmetric-skew symmetric decomposition and the polar description of complex numbers is analogous to the polar decomposition.

$$
\begin{array}{lll}
z=a+b i, & \Rightarrow & A=S+K \\
z=|z| e^{i \theta} & \Rightarrow & A=R P=P^{\prime} R
\end{array}
$$

This analogy extends in the following ways detailed in the diagram. For complex numbers and square invertible matrices and

$$
z=a+b i=|z| e^{i \theta}, \quad A=U \Sigma V^{T}
$$

- $z \Rightarrow A, z^{*} \Rightarrow A^{T}, z^{-1} \Rightarrow A^{-1}, z^{-*} \Rightarrow A^{-T}$
- $a \Rightarrow S=\frac{1}{2}\left(A+A^{T}\right), b i \Rightarrow K=\frac{1}{2}\left(A-A^{T}\right)$.
- $|z|=\left(z^{*} z\right)^{\frac{1}{2}} \Rightarrow P=\left(A^{T} A\right)^{\frac{1}{2}}=V \Sigma V^{T}, P^{\prime}=\left(A A^{T}\right)^{\frac{1}{2}}=U \Sigma U^{T}$,
- $e^{i \theta} \Rightarrow R=A\left(A^{T} A\right)^{-\frac{1}{2}}=\left(A A^{T}\right)^{-\frac{1}{2}} A=U V^{T}, e^{-i \theta} \Rightarrow R^{T}$,


